

# Biophysics 230: Nuclear Magnetic Resonance Haacke Chapter 11

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# 11: The Continuous and Discrete Fourier Transforms

## The Continuous Fourier Transform

We've already discussed the continuous 1D Fourier transform

$$F(\nu) = \int_{-\infty}^{+\infty} f(x) e^{-i2\pi\nu x} dx$$

and its inverse

$$f(x) = \int_{-\infty}^{+\infty} F(\nu) e^{+i2\pi\nu x} d\nu$$

The 1D signal which arises from the application of magnetic (gradient) fields in the  $x$  direction is

$$s(k) = \int_{-\infty}^{+\infty} \rho(x) e^{-i2\pi kx} dx \quad (11.1)$$

## 11: The Continuous and Discrete Fourier Transforms

If we were able to continuously measure  $s(k)$  for all  $k$ , then we could compute the inverse Fourier transform

$$\rho(x) = \int_{-\infty}^{+\infty} s(k) e^{+i2\pi kx} dk \quad (11.7)$$

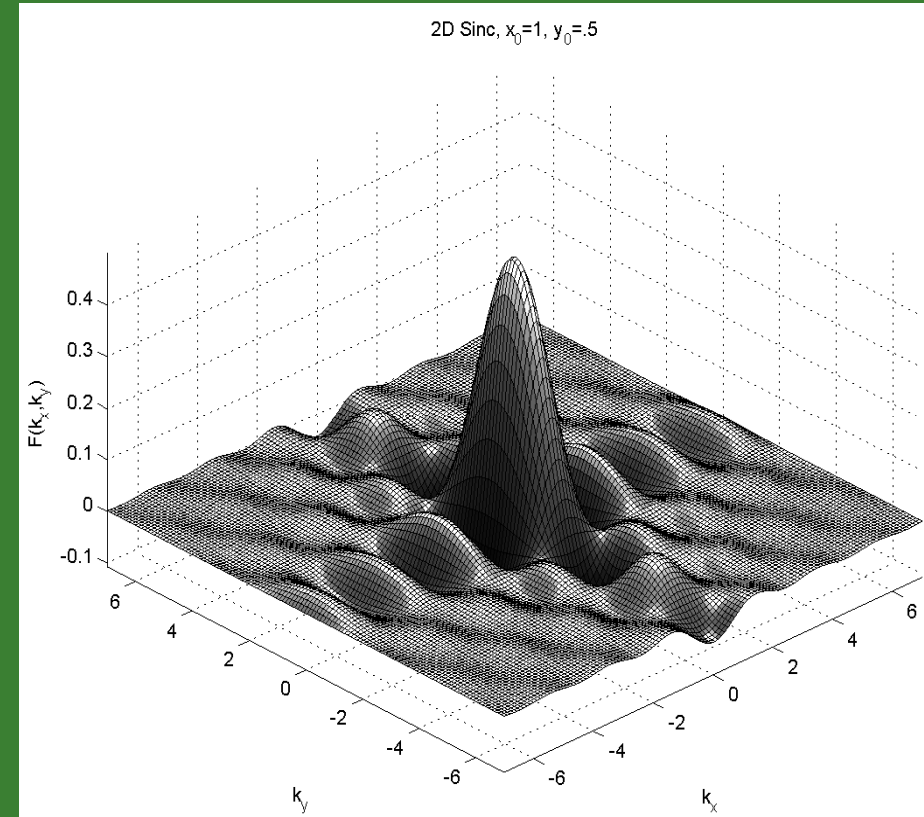
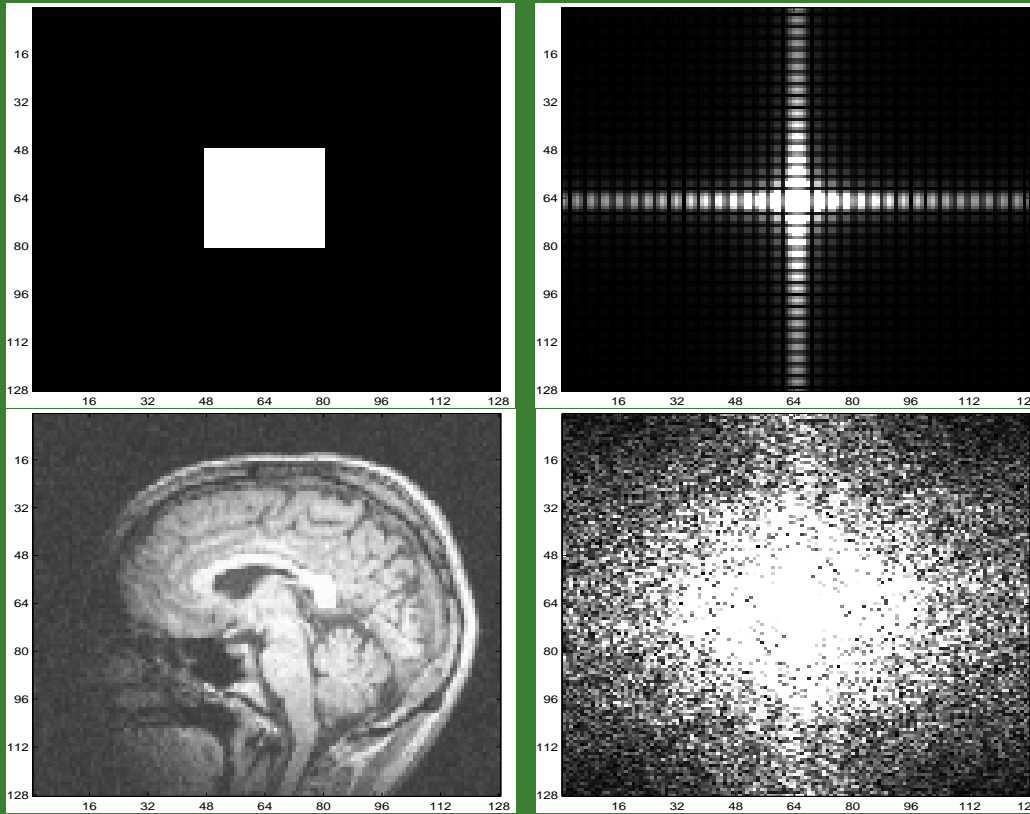
and obtain our spin density.

Unfortunately or fortunately we can only measure/record/observe the signal at discrete time intervals. Every  $\Delta t$ .

If the measured signal is  $s_m(k)$ , the inverse Fourier transform which is the estimated spin density  $\hat{\rho}(x)$  is

$$\hat{\rho}(x) = \int_{-\infty}^{+\infty} s_m(k) e^{+i2\pi kx} dk \quad (11.8)$$

# 11.2: Continuous Transform Properties & Phase Imaging



I can do this too!

Actually the MR image was a .jpg (joint photographic experts group).

## 11.2: Continuous Transform Properties & Phase Imaging

### Complexity of the Reconstructed Image

Our true image  $\rho(x, y)$  is always real,  
our reconstructed image  $\hat{\rho}(x, y)$  is not in general real.

One obvious reason is a constant phase shift.

$$\tilde{s}(k) = e^{i\phi_0} s(k) \quad (11.9)$$

which can arise from the real and imaginary channels being switched or from incorrect demodulation and leads to

$$\hat{\rho}(x) = e^{i\phi_0} \rho(x) \quad (11.10)$$

## 11.2: Continuous Transform Properties & Phase Imaging

### The Shift Theorem

Signal shifted by  $k_0$  in  $k$ -space so truly not  $s(k)$  but truly  $s(k - k_0)$ .

Can happen because of improper demodulation or center of echo at  $k = k_0$  instead of  $k = 0$

The effect on the 1D spin density is

$$\begin{aligned}
 \hat{\rho}(x) &= \int_{-\infty}^{+\infty} s_m(k - k_0) e^{i2\pi kx} dk \\
 &= e^{i2\pi k_0 x} \int_{-\infty}^{+\infty} s_m(k') e^{i2\pi kx} dk' \\
 &= e^{i2\pi k_0 x} \hat{\rho}_{expected}(x)
 \end{aligned} \tag{11.11}$$

in the third line change of variable  $k' = k - k_0$  and  $dk' = dk$ .

The magnitude is not affected,  $|\hat{\rho}(x)| = \hat{\rho}(x)$ .

## 11.2: Continuous Transform Properties & Phase Imaging

There can also be a spatial shift due to an improper read gradient as in Figure 11.2.

### Phase Imaging and Phase Aliasing

Read. Not discussing. I think it is important.

### Duality

The shift theorem also applies to the delta function.

$$\begin{aligned}h(x) &= \delta(x - x_0) \\ H(k) &= e^{-i2\pi kx_0}\end{aligned}\tag{11.14}$$

## 11.2: Continuous Transform Properties & Phase Imaging

### Convolution Theorem (The most important theorems in MRI!)

Get a modified image due to multiplication by function ('filter').

The Fourier transform of a product

$$\mathcal{F} \{g(x) \cdot h(x)\} = G(k) * H(k) \quad (11.16)$$

where

$$G(k) * H(k) \equiv \int G(k') H(k - k') dk' . \quad (11.17)$$

the reverse is also true,

$$\mathcal{F} \{g(x) * h(x)\} = G(k) \cdot H(k)$$

The details of the convolution are learned on Biophysics 03240.

But do know the theorem! There are very important practical implications.

See Figure 11.5.



## 11.2: Continuous Transform Properties & Phase Imaging

### Convolution Associativity

$$a(x) * (b(x) * c(x)) = (a(x) * b(x)) * c(x)$$

### Other Convolution Properties

$$g(x) * h(x) = h(x) * g(x) \quad \text{commutative}$$

$$g(x) * (h_1(x) + h_2(x)) = g(x) * h_1(x) + g(x) * h_2(x) \quad \text{distributive}$$

## 11.2: Continuous Transform Properties & Phase Imaging

**Derivative Theorem** (We already discussed this)

$$\begin{aligned}
 \mathcal{F}\{f'(x)\} &= \mathcal{F}\left\{\frac{df(x)}{dx}\right\} \\
 &= \int_{-\infty}^{+\infty} \left[\frac{df(x)}{dx}\right] e^{-i2\pi kx} dx \\
 &= \left[ f(x)e^{-i2\pi kx} \Big|_{-\infty}^{+\infty} - (-i2\pi k) \int_{-\infty}^{+\infty} f(x)e^{-i2\pi kx} dx \right] \\
 &= i2\pi k F(k)
 \end{aligned} \tag{11.22}$$

in the second to last line we integrated by parts

$$\int g(x)h'(x) dx = g(x)h(x) - \int g'(x)h(x)dx$$

where  $g(x) = e^{-i2\pi kx}$  and  $h'(x) = \frac{df(x)}{dx}$ .

FT of derivative (1D image)  $f'(x)$  by multiplying  $F(k)$  by  $i2\pi k$ .

## 11.2: Continuous Transform Properties & Phase Imaging

### Example:

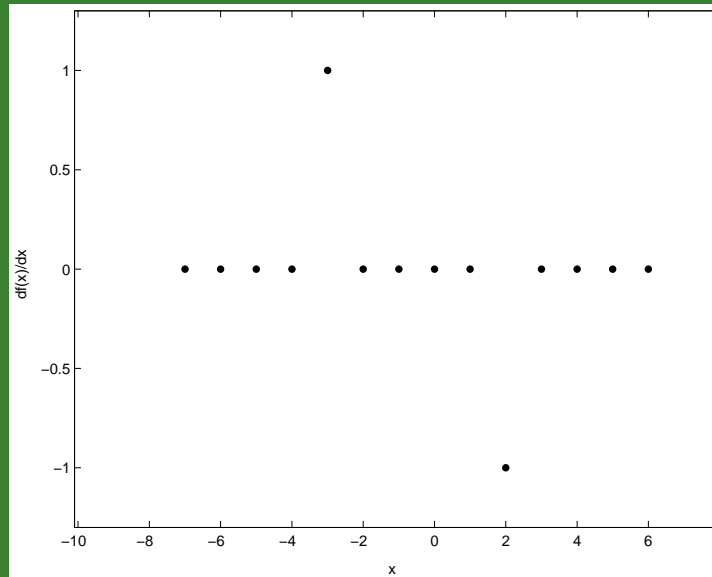
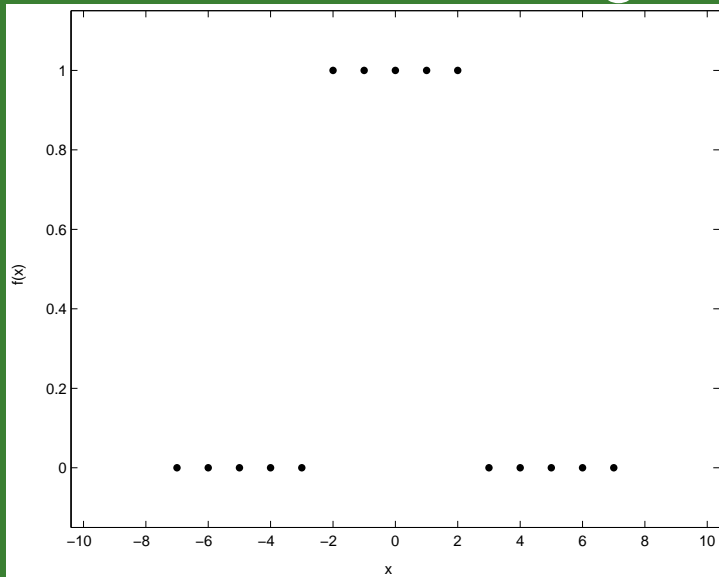
When taking the derivative of an image, it is actually discrete differences.

Consider the following discrete function  $f(x)$ .

Define the derivative at a point  $x_i$  to be

$$f'(x_i) = \frac{f(x_i) - f(x_{i-1})}{x_i - x_{i-1}}.$$

Then the function along with its derivative are



This is also true in two dimensions.

## 11.2: Continuous Transform Properties & Phase Imaging

Now look at Figure 11.6 in the book.

A 1D derivative of the magnitude image is taken in the vertical direction.

The magnitude means derivative is strictly positive.

We could also take a 2D derivative, the gradient.

The derivative or gradient tells us where the edges are in the image.

Can also do the Laplacian.

What about integrals?

### Fourier Transform Symmetries

Read.

## 11.2: Continuous Transform Properties & Phase Imaging

### Summary of 1D Fourier Transform Properties

| Property          | Function                      | Transform                                |
|-------------------|-------------------------------|--|
| <b>Linearity</b>  | $af(x) + bg(x)$               | $aF(k) + bG(k)$                          |
| <b>Similarity</b> | $f(ax)$                       | $\frac{1}{ a }F\left(\frac{k}{a}\right)$ |
| <b>Shifting</b>   | $f(x - a)$                    | $e^{-i2\pi ka}F(k)$                      |
| <b>Derivative</b> | $\frac{d^\ell f(x)}{dx^\ell}$ | $(i2\pi k)^\ell F(k)$                    |

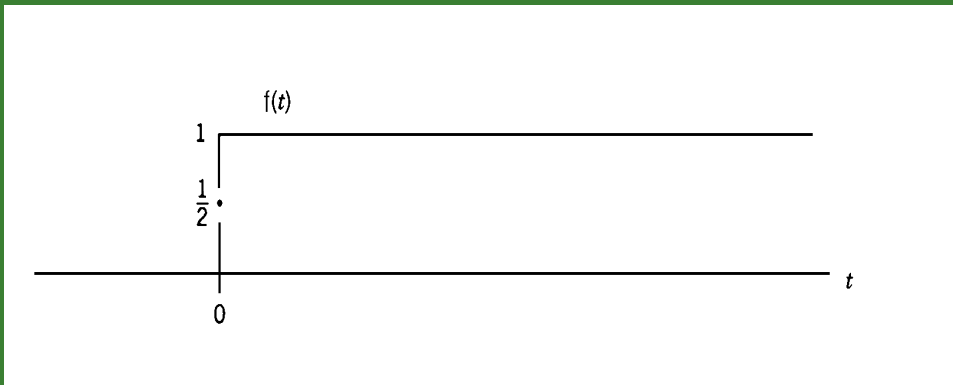
## 11.3: Fourier Transform Pairs

### Heaviside Function

The Heaviside function is

$$\Theta(k) = \begin{cases} 1 & k > 0 \\ \frac{1}{2} & k = 0 \\ 0 & k < 0 \end{cases} . \quad (11.24)$$

Here is a picture of the Heaviside function. Often it doesn't have the  $\frac{1}{2}$ .



The inverse Fourier transform of the Heaviside function can be found as

$$h_{\Theta}(x) = \mathcal{F}^{-1}\{\Theta(k)\}$$

## 11.3: Fourier Transform Pairs

$$\begin{aligned}
 h_{\Theta}(x) &= \lim_{\epsilon \rightarrow 0^+} \int_{-\infty}^{+\infty} \Theta(k) e^{i2\pi kx} e^{-2\pi\epsilon|k|} dk \\
 &= \lim_{\epsilon \rightarrow 0^+} \left[ \int_{-\infty}^{0^-} \overbrace{\Theta(k)}^{=0} e^{i2\pi kx} e^{-2\pi\epsilon(-k)} dk \right. \\
 &\quad \left. + \int_0^{+\infty} \Theta(k) e^{i2\pi kx} e^{-2\pi\epsilon k} dk \right] \\
 &= \lim_{\epsilon \rightarrow 0^+} \int_0^{+\infty} \Theta(k) e^{-2\pi(\epsilon - ix)k} dk \\
 &= \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi(\epsilon - ix)} e^{-2\pi(\epsilon - ix)k} \Bigg|_0^{+\infty}
 \end{aligned}$$

\* $e^{-2\pi|k|x}$  to eliminate convergence ambiguities

## 11.3: Fourier Transform Pairs

$$\begin{aligned}h_{\Theta}(x) &= \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi(\epsilon - ix)} \\&= \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi} \frac{1}{(\epsilon - ix)} \frac{(\epsilon + ix)}{(\epsilon + ix)} \\&= \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi} \frac{(\epsilon + ix)}{(\epsilon^2 + x^2)} \\&= \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi} \left[ \frac{\epsilon}{(\epsilon^2 + x^2)} + \frac{ix}{(\epsilon^2 + x^2)} \right] \\&= \frac{1}{2\pi} \delta(x) + \frac{i}{2\pi} P\left(\frac{1}{x}\right)\end{aligned}\tag{11.25}$$



## 11.3: Fourier Transform Pairs

### Lorentzian Form

The two sided (double exponential)

$$f(k) = e^{-a|k|}$$

inverse Fourier transform of this is

$$\begin{aligned}
 F(x) &= \int_{-\infty}^{+\infty} e^{-a|k|} e^{i2\pi kx} dk \\
 &= \int_{-\infty}^{+\infty} e^{-a|k|} e^{i2\pi kx} dk \\
 &= \int_{-\infty}^0 e^{(i2\pi x + a)k} dk + \int_0^{+\infty} e^{(i2\pi x - a)k} dk \\
 &= \frac{1}{(i2\pi x + a)} e^{(i2\pi x + a)k} \Big|_{-\infty}^0 + \frac{1}{(i2\pi x - a)} e^{(i2\pi x - a)k} \Big|_0^{+\infty}
 \end{aligned} \tag{11.26}$$

## 11.3: Fourier Transform Pairs

### Lorentzian Form

$$\begin{aligned}
 F(x) &= \frac{1}{(i2\pi x + a)} - \frac{1}{i2\pi x - a} \\
 &= \frac{1}{(i2\pi x + a)(i2\pi x + a)} - \frac{1}{(i2\pi x - a)(-i2\pi x + a)} \\
 &= \frac{i2\pi x + a}{(2\pi x)^2 + a^2} + \frac{-i2\pi x + a}{(2\pi x)^2 + a^2} \\
 &= \frac{2a}{(2\pi x)^2 + a^2}
 \end{aligned}$$

which is called “Lorentzian” form.

## 11.3: Fourier Transform Pairs

### The Sampling Function

The sampling or comb function  $u(k)$  is defined to be the sum of a (doubly) infinite number of delta functions each  $\Delta k$  apart.

$$u(k) = \Delta k \sum_{p=-\infty}^{+\infty} \delta(k - p\Delta k) \quad (11.28)$$

When we measure/record/observe/sample the signal  $s(k)$

at discrete  $k$ -space (time) points  $\Delta k$  ( $\Delta t$ ) apart,

this is equivalent to multiplying the continuous signal by the comb function.

## 11.4 The Discrete Fourier Transforms

The inverse Fourier transform of the comb function is

$$\begin{aligned}
 U(x) &= \mathcal{F}\{u(k)\} \\
 &= \Delta k \sum_{p=-\infty}^{+\infty} \int \delta(k - p\Delta k) e^{i2\pi kx} dk \\
 &= \Delta k \sum_{p=-\infty}^{+\infty} e^{i2\pi p\Delta kx}
 \end{aligned} \tag{11.29}$$

where

$$\sum_{n=-\infty}^{+\infty} e^{i2\pi na} = \sum_{m=-\infty}^{+\infty} \delta(a - m) \tag{11.30}$$

and thus

$$U(x) = \sum_{q=-\infty}^{+\infty} \delta\left(x - \frac{q}{\Delta k}\right) \tag{11.31}$$

## 11.4 The Discrete Fourier Transforms

The DFT is an approximation to the continuous FT.

Often in theory assume  $s_m(k)$  is continuous so we could perform an IFT.

In reality we sample the signal  $s_m(k)$  at discrete  $k$ -space (time) points.

With a length scale  $L$ , the discrete Fourier transform

$$G\left(\frac{p}{L}\right) \equiv \mathcal{D}(g) = \sum_{q=-n}^{n-1} g\left(\frac{qL}{2n}\right) e^{-\frac{i2\pi pq}{2n}} \quad (11.32)$$

$\Delta k = 1/L$  and  $L = 2n\Delta x$

$$G(p\Delta k) = \sum_{q=-n}^{n-1} g(q\Delta x) e^{-\frac{i2\pi pq\Delta x\Delta k}{2n}} \quad (11.33)$$

$$g\left(\frac{ql}{2n}\right) \equiv \mathcal{D}^{-1}(G) = \frac{1}{2n} \sum_{p=-n}^{n-1} G\left(\frac{p}{L}\right) e^{\frac{i2\pi pq}{2n}} g(q\Delta x) \quad (11.34)$$

$$= \frac{1}{2n} \sum_{p=-n}^{n-1} G(p\Delta k) e^{\frac{i2\pi pq}{2n}} \quad (11.35)$$

## 11.4 The Discrete Fourier Transforms

### Other DFT Pair Parameterizations

$$F(u) = \frac{1}{N} \sum_{x=0}^{N-1} f(x) e^{-i2\pi ux/N}$$

$$f(x) = \sum_{u=0}^{N-1} F(u) e^{+i2\pi ux/N}$$

$$F(u) = \sum_{x=1}^N f(x) e^{-i2\pi(u-1)(x-1)/N}$$

$$f(x) = \frac{1}{N} \sum_{u=1}^N F(u) e^{+i2\pi(u-1)(x-1)/N}$$

\* Matlab's "fft" and "ifft" use the second one.

## 11.5 Discrete Transform Properties

### The Discrete Convolution Theorem

The convolution theorem holds for the discrete convolution just as for the continuous convolution. Using a shorthand notation,

$$g_1(q)g_2(q) \stackrel{\mathcal{D}}{=} G_1(p) * G_2(p) \quad (11.39)$$

and

$$G_1(p) * G_2(p) = \frac{1}{2n} \sum_{r=-n}^{n-1} G_1(r)G_2(p-r) \quad (11.40)$$

also

$$g_1(q) * g_2(q) \stackrel{\mathcal{D}}{=} G_1(p)G_2(p) \quad (11.41)$$

## 11.5 Discrete Transform Properties

### Summary of Discrete Fourier Transform Properties

Table 11.4 summarizes the DFT properties.